

Physically-Based High-Level System Model of a MEMS-Gyroscope for the Efficient Design of Control Algorithms

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ABSTRACT

We present a high-level model of a dual gimballed mass gyroscope, which provides an accurate physical description of the impact of external and internal disturbances on the output signal, but which also allows for the efficient analysis and optimization of the full sensor system including the electronic circuitry for drive, control, and signal conditioning. The dynamics of the gyroscope is described by a reduced-order high-level model. Internal disturbances (manufacturing tolerances, e.g.) as well as external impact factors (pressure-dependent viscous damping, shock, vibrations of the housing, etc.) are included in the model equations by introducing physically-based, parameterized functions extracted from detailed FEM or mixed-level simulations. Exemplary simulations investigating the impact of package vibrations on the sensor output signal prove the efficiency of our model and the practicality of our approach.

1 INTRODUCTION

Robust sensor systems designed for applications in harsh environments (e.g., automobiles) have to be equipped with additional “smart functionalities” to compensate environmental impacts (e.g., ambient pressure variations or package vibrations) or internal imperfections due to manufacturing tolerances.

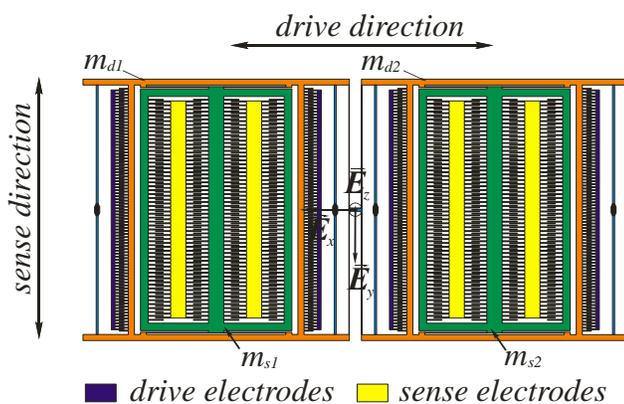


Fig. 1: Structure of the dual gimballed mass gyroscope.

In order to ensure a proper and reliable sensor operation, adequate concepts for sensor control and signal conditioning must be developed, which are based on a profound understanding of the effects disturbing the sensor output signal. To this end, we derived a high-level model of the dual gimballed mass gyroscope depicted in Fig. 1, which provides an accurate physical description of the impact of external and internal disturbances on the output signal, but, at the same time, allows for the efficient analysis and optimization of the full sensor system including the electronic circuitry for drive, control, and signal conditioning and, thus, is perfectly suited for the co-design of the transducer elements and the control electronics

2 HIGH-LEVEL GYROSCOPE MODEL

2.1 Modeling Approach

As it would be much too time- and memory-consuming to perform a complete continuous-field FEM analysis of the gyroscope including all possible mechanical configurations, we focus on the most relevant modes of motion and decompose the device structure into four rigidly moving

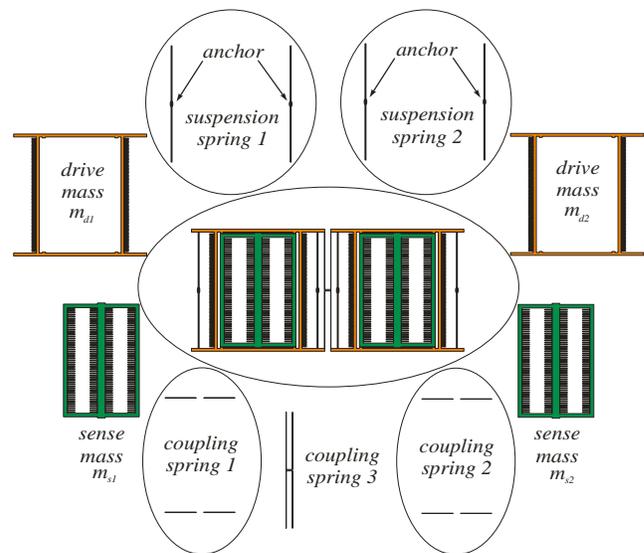


Fig. 2: Decomposition of the gyroscope into four rigidly moving bodies connected by mass-less springs.

mechanical substructures, which are hinged and inter linked by springs (Fig. 2). The dynamics of this simplified structure is described by the Langrangian equations of motion, leading to a reduced-order model of the complete gyroscope with 24 degrees of freedom [1]:

$$\begin{aligned} \mathbf{M} \cdot \ddot{\mathbf{u}} + (\mathbf{C} + \mathbf{C}_\Omega) \cdot \dot{\mathbf{u}} + (\mathbf{K} + \mathbf{K}_\Omega) \cdot \mathbf{u} &= \\ &= -\mathbf{M} \cdot \underline{a}_o + \underline{F}_{el} \end{aligned} \quad (1)$$

Here \mathbf{K} and \mathbf{M} denote the stiffness and the mass matrices, \mathbf{C} the damping matrix, which accounts for viscous friction due to the surrounding air, \underline{F}_{el} the electrostatic forces actuating the drive masses, and \underline{a}_o the external linear acceleration acting on the substrate frame. The measurand, the angular velocity $\vec{\Omega}$, enters the system through the terms \mathbf{C}_Ω and \mathbf{K}_Ω , which comprise the contributions of the Coriolis and the centrifugal force to the stiffness and the damping matrix, respectively.

2.2 Mechanical Model Parameters

The mechanical model parameters like the mass matrix \mathbf{M} and the stiffness matrix \mathbf{K} were extracted from detailed 3D FEM simulations by calculating the respective inertial forces and spring constants for each rotational and translational degree of freedom. The three fundamental resonance frequencies and the associated mode shapes of the gyroscope as obtained from the harmonic solutions of the homogeneous equation system (1) (i.e. $\underline{a}_o = 0$, $\underline{F}_{el} = 0$) conform well with those calculated by modal FEM analysis (see Fig. 3 and Tab. 1).

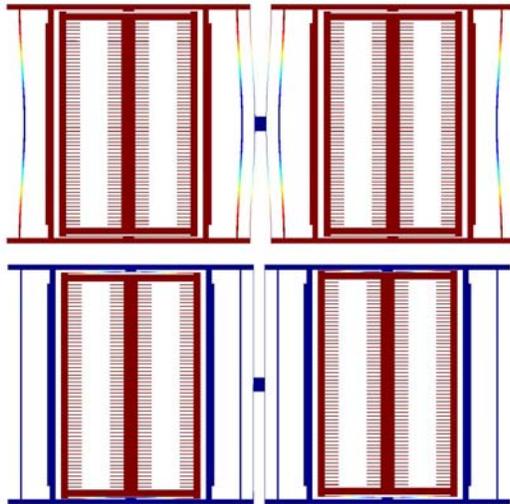


Fig. 3: Results of FEM modal analysis of the gyroscope structure.

Top: anti-phase drive mode at 10.654 kHz
Bottom: sense mode at 10.527 kHz

Resonance frequencies of the gyroscope	FEM	harmonic solution
in-phase drive mode	10.301 kHz	10.459 kHz
sense mode	10.527 kHz	10.677 kHz
anti-phase drive mode	10.654 kHz	10.723 kHz

Table 1: Three fundamental mechanical resonance frequencies of the gyroscope.

It shows that manufacturing fluctuations such as, e.g., etch bias tolerances have a considerable impact on the mass and stiffness parameters.

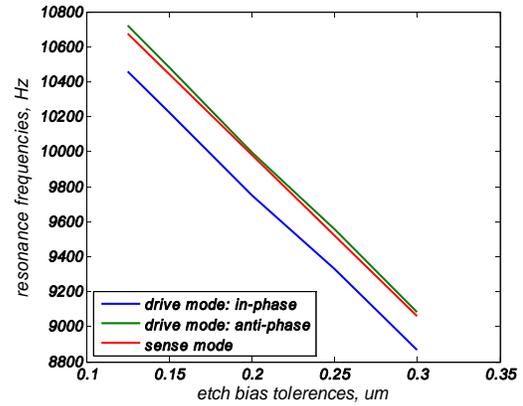


Fig. 4: Impact of manufacturing tolerances on the resonance frequencies of the drive and the sense masses.

The dependencies from these parameter variations have been analyzed (see Fig. 4) and extracted by detailed FEM simulations and, subsequently, fed into the high-level model as parametrized functions.

2.3 Viscous Damping Model

Viscous damping due to the surrounding air is, by its nature, a distributed effect, which cannot be properly described by analytical compact models and which is also difficult to extract from 3D coupled-domain FEM simulations, especially for geometrically complex microstructures. Therefore, we included this effect by introducing modal quality factors \hat{C}_i for all relevant modes of motion:

$$\hat{C}_i = \frac{\sqrt{\hat{K}_i \cdot \hat{M}_i}}{\hat{Q}_i} \quad (2)$$

where \hat{K}_i , \hat{M}_i and \hat{Q}_i denote the stiffness, mass and quality factors of the i -th eigenmode, respectively. The \hat{C}_i are then converted into the spatial damping matrix \mathbf{C} by applying modal transformation techniques.

Since it is computationally quite expensive and for the required accuracy of the system model also not necessary to determine the quality factors for a large number of eigenmodes, we focussed on the two most relevant modes of motion, i.e. the sense mode, which is mainly affected by squeeze film damping, and the drive mode, whose dynamics is dominated by slide film damping. In detail, the modal quality factor of the drive mode can be determined by:

$$\frac{1}{Q_{drive}} = \frac{1}{Q_{sl_fr}} + \frac{1}{Q_{comb}} \quad (3)$$

where Q_{sl_fr} represents the quality factor of the drive and sense frames and Q_{comb} the quality factor of the comb drives. For the calculation of both quality factors we apply the following expression for slide film damping given in [2]:

$$Q_i = \frac{\hat{M}_i \cdot \omega_i \cdot h}{\eta_{eff} \cdot A_{eff}} \quad (4)$$

Here ω_i stands for the angular eigenfrequency of the i -th eigenmode (drive frequency), h for the distance between the moving frame and the substrate or, alternatively, the distance between the comb fingers, A_{eff} for the relevant effective area, and η_{eff} is the effective viscosity accounting for gas rarefaction effects in the low pressure regime and/or small structural dimensions [2]:

$$\eta_{eff} = \frac{\eta_T}{1 + 2 \cdot K_n + 0.2 \cdot K_n^{0.788} e^{-K_n/10}} \quad (5)$$

(with η_T = viscosity of air under normal pressure conditions at room temperature and K_n = Knudsen number).

The modal quality factor of the sense mode is dominated by squeeze film damping and was determined by applying the mixed-level modeling approach proposed in [3]. This approach takes advantage from the small Reynolds numbers and the large aspect ratios typically encountered in MEMS structures and reduces the degree of complexity by replacing the non-linear and highly complicated Navier-Stokes equation by the well-known Reynolds' equation [4]:

$$\nabla \left(\frac{\rho h^3}{12 \eta_{eff}} \nabla p(x, y) \right) = \frac{\partial}{\partial t} (\rho h) \quad (6)$$

The pressure distribution $p(x, y)$ underneath the moving plates is calculated by discretizing this equation to form a fluidic Kirchhoffian network and solving it by the use of a

standard circuit simulator. Edge effects and perforations in the structure are taken into account by introducing physically-based compact models at the respective locations. The pressure-dependent quality factors of the sense mode, as they have been extracted from these mixed-level simulations and incorporated in the high-level model, are shown in Fig. 5.

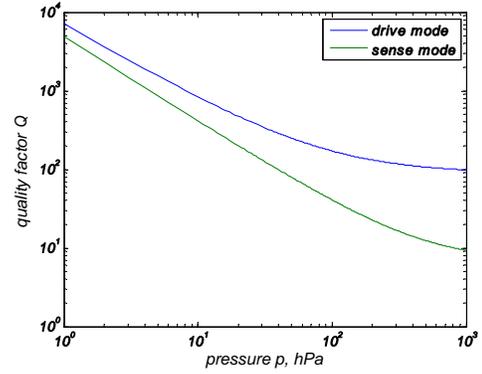


Fig. 5: Pressure-dependent quality factors of the drive and the sense mode of the gyroscope.

3 IMPACT OF ENVIRONMENTAL DISTURBANCES ON THE SENSOR SYSTEM

By their nature, inertial sensors are sensitive to any kind of mechanical forces originating from their environment. Thus, their operation is always affected by disturbing effects arising under real-world operating conditions like shock, cross-talk between sensing and non-sensing axes, and vibrations of the housing. These impact factors enter our sensor model in a physically transparent manner through the parameters C_{Ω} and K_{Ω} and the right-hand side of the model equations, respectively; hence, it provides the proper basis for studying their influence on the sensor system on the whole and, therefore, it is well suited for the design of adaptive control algorithms, which minimize or even eliminate the environmental disturbances.

As an example, we demonstrate the efficiency of our high-level model by analyzing the impact of housing vibrations on the transient sensor response. To this end, we first consider the package alone and extract its eigenmodes and resonance frequencies from FEM calculations [1]. Tab. 2 shows the results for two exemplary standard package types (PSOIC and QFP).

		PSOIC	QFP
in-plane	x-direction	14.654 kHz	30.379 kHz
in-plane	y-direction	14.835 kHz	30.387 kHz
out-of-plane	z-direction	25.666 kHz	43.656 kHz

Table 2: Vibration modes of the PSOIC and QFP packages

