



# PSP-based scalable compact FinFET model

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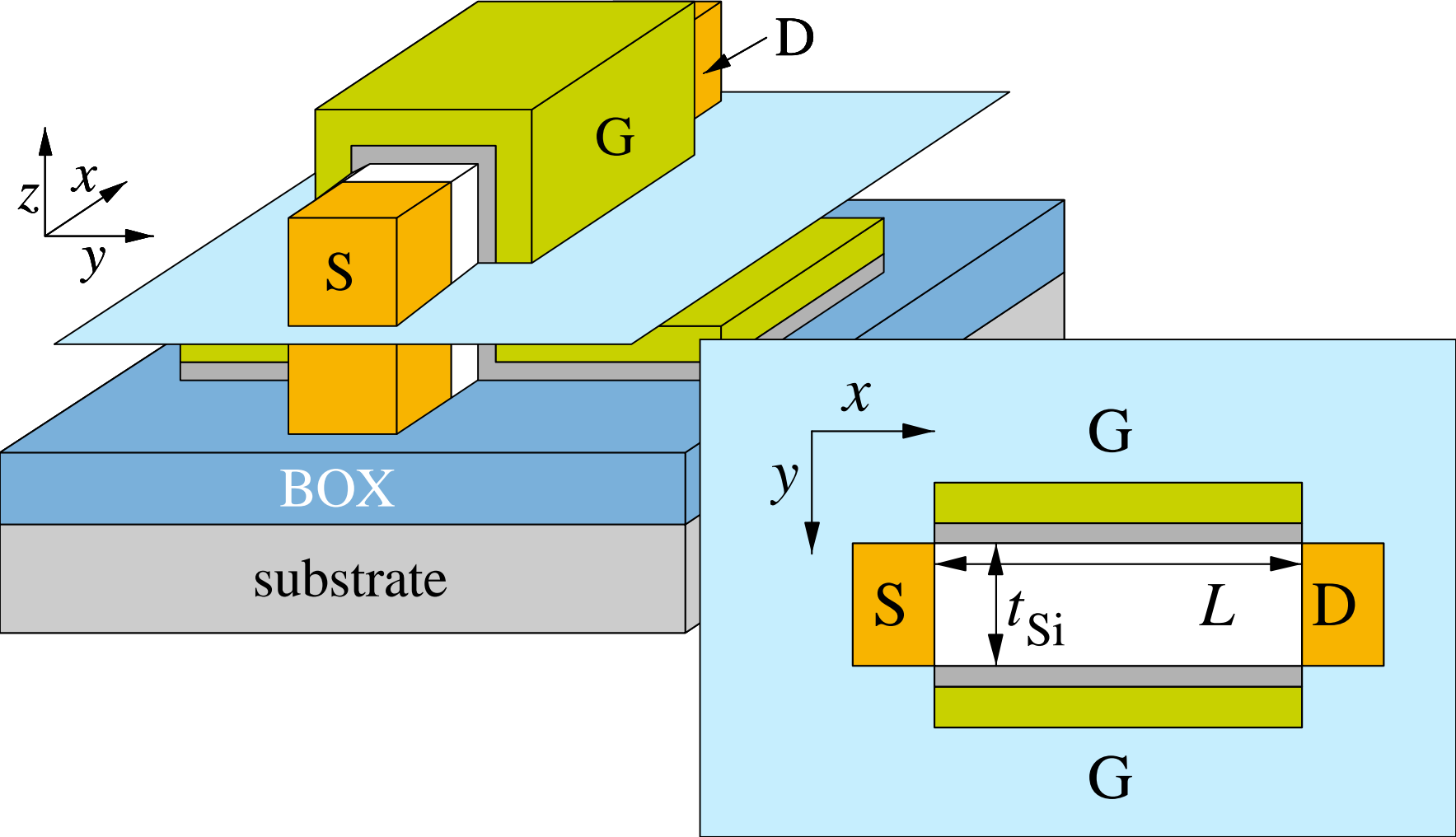


# introduction

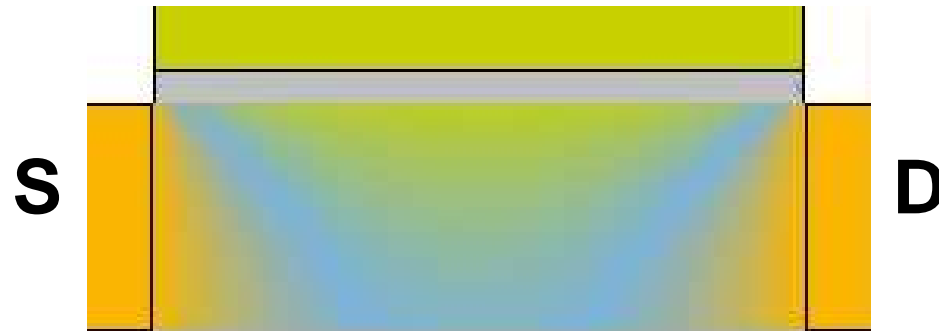
- ▶ FinFETs are candidates to replace bulk MOSFETs
  - good control of short channel effects
- ▶ compact model development
  - compact model required for circuit design
  - no (public domain) FinFET models available
- ▶ extending the PSP model-family
  - bulk CMOS
  - PD-SOI
  - FD-SOI
  - FinFET
  - ...
- ▶ FinFET is fully depleted device
  - fundamentally different electrostatics
  - touches the heart of a compact model



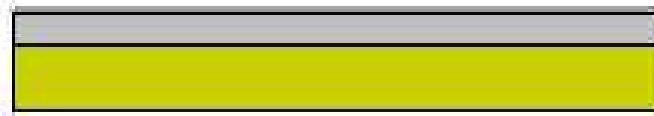
# device



# channel control in FinFET



(animation)



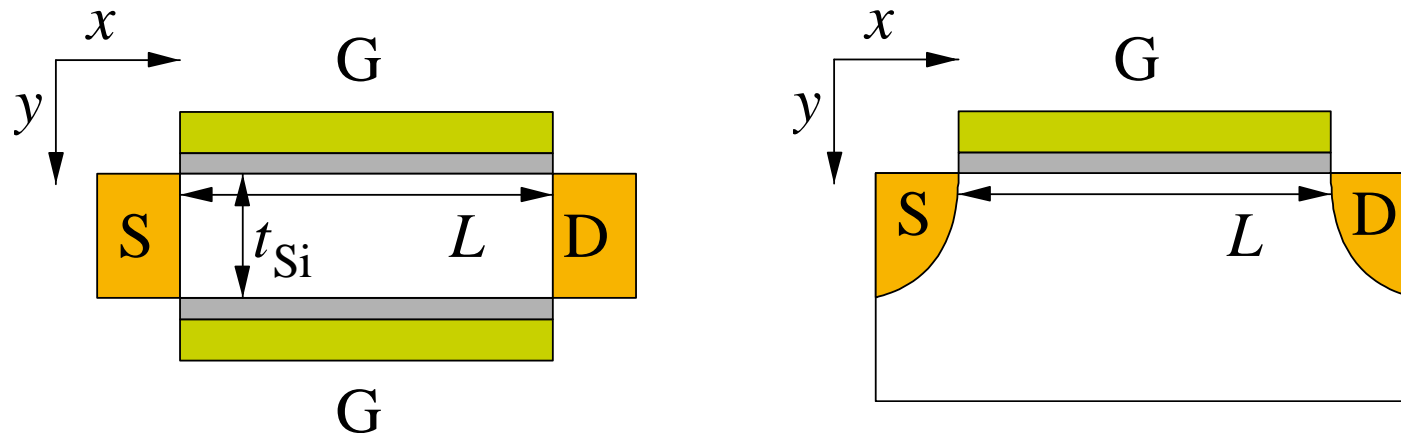
# outline

- ▶ introduction
- ▶ core model
  - surface potential equation & its solution
  - currents and charges
  - new model
- ▶ geometry scaling
  - parameter extraction
  - results
- ▶ conclusion



**core model**

# surface potential equation (i)



- ▶ 2D Poisson equation 
$$-\left(\frac{\partial^2 \psi}{\partial x^2} + \frac{\partial^2 \psi}{\partial y^2}\right) = \frac{\rho}{\epsilon}$$

- ▶ gradual channel approximation

- ▶ 1D Poisson equation 
$$-\frac{\partial^2 \psi}{\partial y^2} = \frac{\rho}{\epsilon}$$

- ▶ charge density 
$$\rho = -q \cdot N_A \cdot \underbrace{\left(1 + e^{-q \cdot (V + \phi_B) / k \cdot T} \cdot e^{q \cdot \psi / k \cdot T}\right)}_{\text{electrons}} \underbrace{- e^{-q \cdot \psi / k \cdot T}}_{\text{holes}}$$

## surface potential equation (ii)

▶ normalize  $\frac{\partial^2 \psi}{\partial y^2} = f(\psi)$   $f(\psi) = (1 + e^{\psi-v} - e^{-\psi})/2$

▶ integrate once  $(\psi')^2 - F(\psi) = \alpha$

$$F(\psi) = 2 \int f(\psi) d\psi = \psi + e^{-v} (e^{\psi} - 1) + (e^{-\psi} - 1)$$

$\alpha \rightarrow$  integration constant

# surface potential equation (iii) – bulk MOSFET

$$(\psi')^2 - F(\psi) = \alpha$$

▶ BC1:  $\psi' = C_{ox} (\psi_s - V_{GB}^*)$

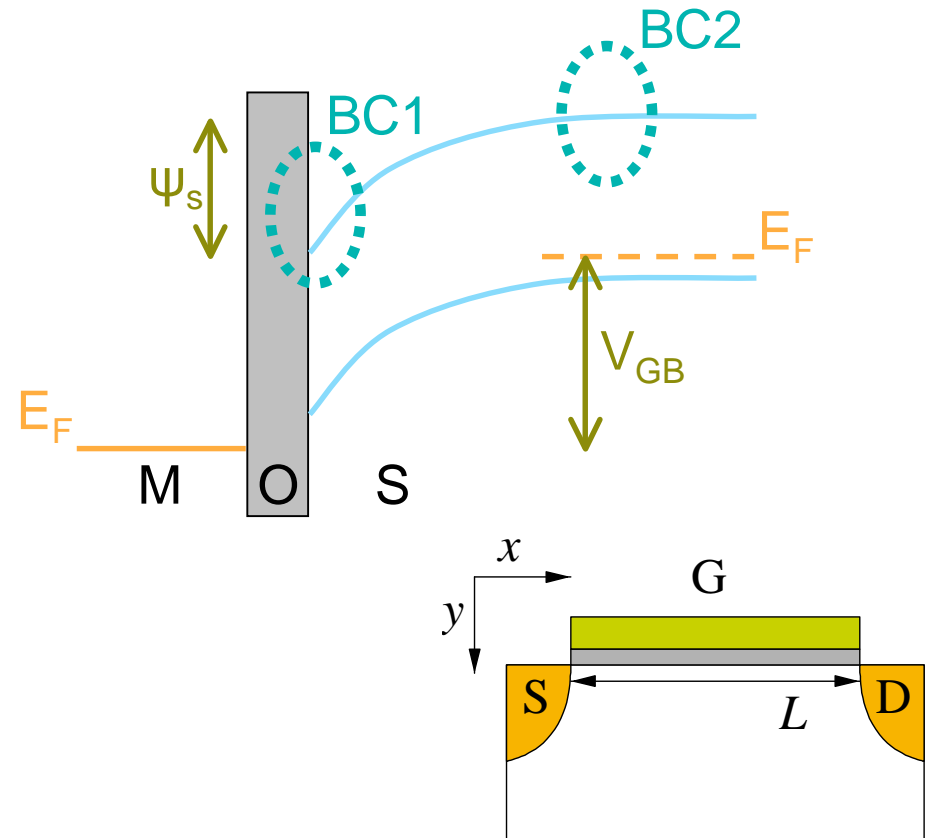
▶ BC2:  $\psi' = 0$

▶  $\alpha = 0$

▶  $\psi_s$  can be solved from:

$$(\psi_s - V_{GB}^*)^2 = F(\psi_s)$$

▶ note:  $\psi(y)$  is *not* yet known



# surface potential equation (iv) – FinFET

$$(\psi')^2 - F(\psi) = \alpha$$

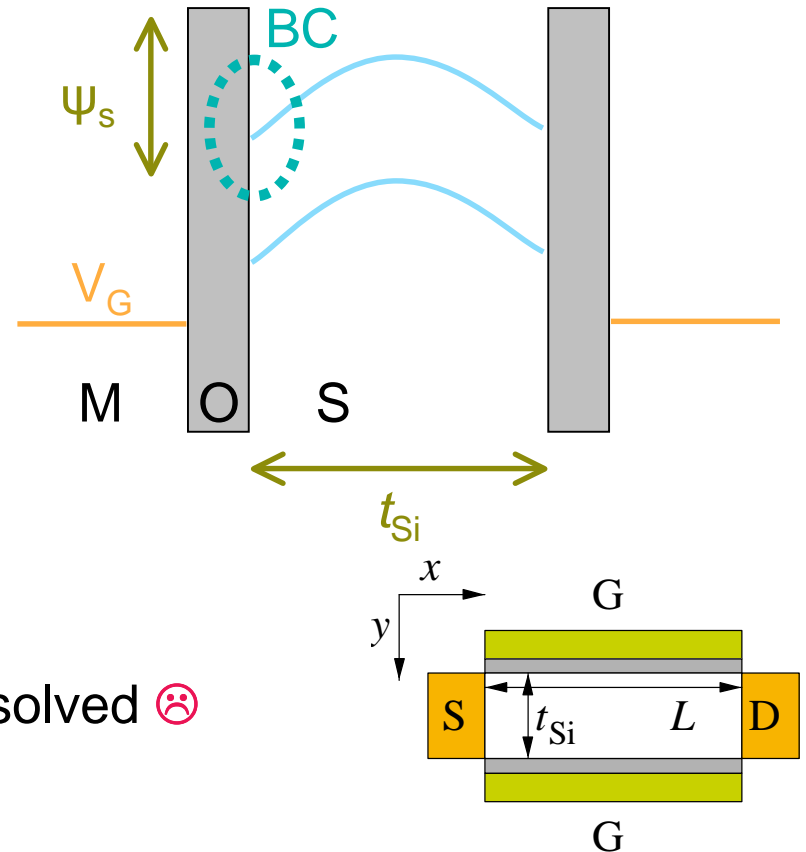
▶ BC:  $\psi' = C_{ox} \cdot (\psi_s - V_G^*)$

▶  $\psi_s$  cannot yet be solved

▶ need second integration step

$$t_{Si} = \int_{\psi_{s,1}}^{\psi_{s,2}} \frac{d\psi}{\sqrt{F(\psi) + \alpha}}$$

▶ two equations (including integral) to be solved ☹



# surface potential equation (v)

- ▶ approximation required
- ▶ 1. neglect holes
  - accumulation region not important for normal MOS operation
  - floating body → holes only visible at very low frequency
- ▶ 2. neglect doping
  - doping typically low →  $t_{\text{Si}} \ll L_D$
  - doping hardly influences potential
- ▶ now integral can be performed

$$\psi(\mathbf{x}, y) = v(\mathbf{x}) - \ln \left[ \frac{\cos^2(\sqrt{|\alpha(\mathbf{x})|} y / 2)}{|\alpha(\mathbf{x})|} \right]$$

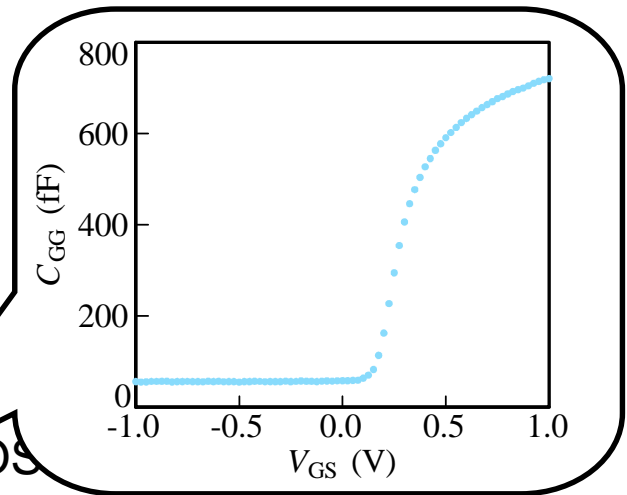
- ▶ use boundary conditions → single algebraic equation
- ▶  $\psi_s$  can be solved (implicit equation)

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# outline

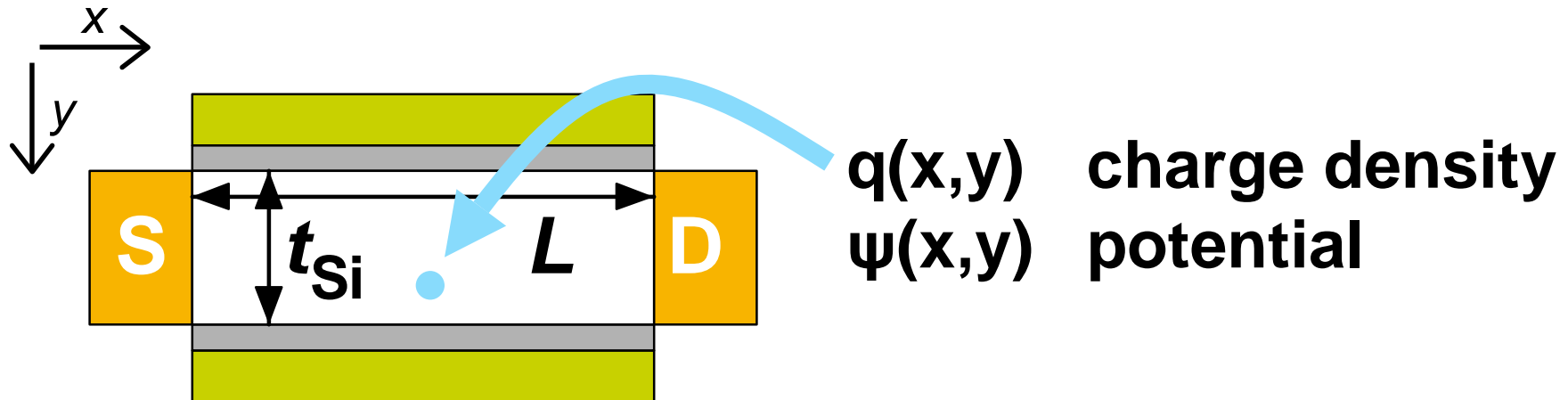
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# currents and charges (i)

- ▶ drift-diffusion gives current density  $J_D(y)$  at any *point* in channel

$$-\frac{J_D(y)}{\mu \cdot h} = q(x, y) \cdot \frac{dv}{dx}(x) = \underbrace{q(x, y) \cdot \frac{d\psi}{dx}(x, y)}_{\text{drift}} - \underbrace{\frac{dq}{dx}(x, y)}_{\text{diffusion}}$$



## currents and charges (ii)

- ▶ drain current (Pao-Sah)

$$I_D = -\mu \cdot \frac{h}{L} \int_0^L \int_{-t_{si}/2}^{t_{si}/2} q(x, y) \cdot \frac{dv}{dx}(x) dy dx$$

- ▶ gate charge

$$Q_G = -h \int_0^L \int_{-t_{si}/2}^{t_{si}/2} q(x, y) dy dx$$

- ▶ integrals can be performed exactly [1]

[1] H. Lu and Y. Taur, TED 53, p. 1126 (2006)

## currents and charges (iii)

- ▶ exact result for gate charge:

$$Q_G \propto \frac{1}{I_D} \left( \left[ \theta^2 - \ln(\cos^2 \theta) - 2\theta \cdot \tan \theta + \frac{4}{3c_{ox} t_{Si}} \theta^3 \tan^3 \theta + h(\theta) \right]_{\theta_0}^{\theta_L} \right)$$

$$h(\theta) = -\frac{2i}{3} \theta^3 - \ln(2) \cdot \theta^2 + (1 - \theta^2) \cdot \ln(\cos \theta) \\ - \frac{\theta^2}{2 \cos^2 \theta} + \theta \tan \theta + i\theta \cdot \text{Li}_2(-e^{2i\theta}) - \frac{1}{2} \text{Li}_3(-e^{2i\theta})$$

$$\theta = \frac{t_{Si}}{4} \sqrt{|\alpha|}$$

$\text{Li}_n(x) \rightarrow$  polylogarithmic function

# current and charges (iv)

- ▶ disadvantages of exact solution
  - complicated expressions (involving complex numbers)
  - special functions (polylogarithms)
    - not available in standard (C-)libraries
    - computationally inefficient
  - difficult to include SCEs
- ▶ not ideal for compact model

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  - [new model](#)
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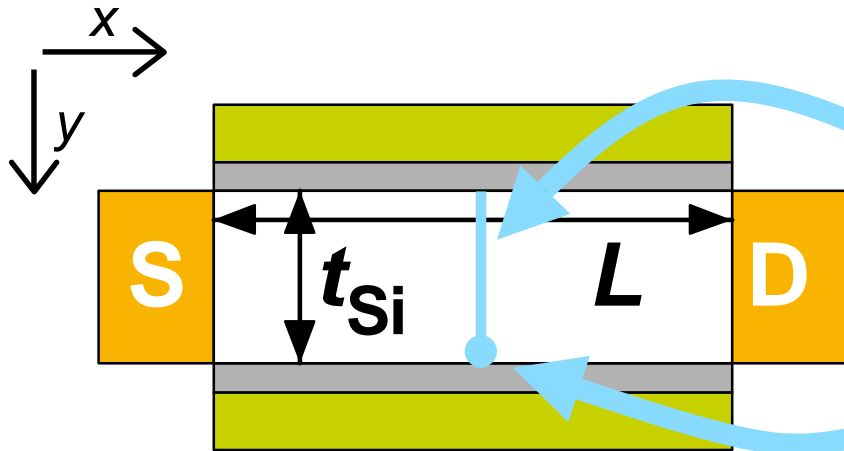


# new model – derivation

$$\frac{-J_D(y)}{\mu \cdot h} = \underbrace{q(x, y)} \cdot \frac{dv}{dx}(x) = \underbrace{q(x, y)} \cdot \frac{d\psi}{dx}(x, y) - \underbrace{\frac{dq}{dx}(x, y)}$$

integrate over y

$$\frac{-I_D}{\mu \cdot h} = Q(x) \cdot \frac{dv}{dx}(x) = \tilde{Q}(x) \cdot \frac{d\psi_s}{dx}(x) - \frac{dQ}{dx}(x)$$



$Q(x)$  (integrated) charge density

$\psi_s(x)$  surface potential

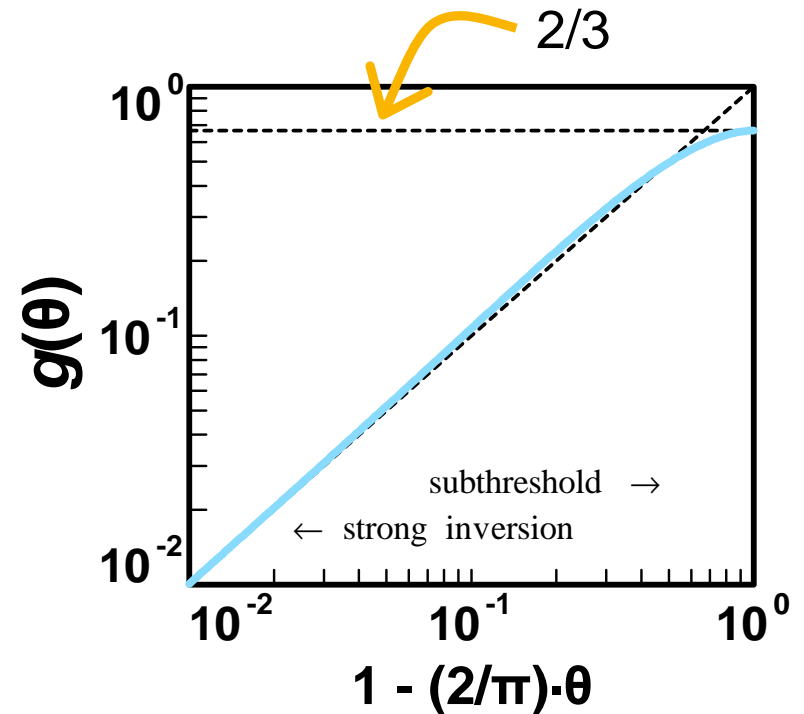
# new model – derivation

▶  $\tilde{Q} = \int_{-t_{Si/2}}^{t_{Si/2}} q \cdot \frac{d\psi}{dx} dy \Big/ \frac{d\psi_s}{dx}$  is related to  $Q$ , but they are *not* equal

▶ in fact:  $\tilde{Q} = Q \cdot \left[ 1 + \frac{\gamma}{4} \cdot g(\theta) \right]$

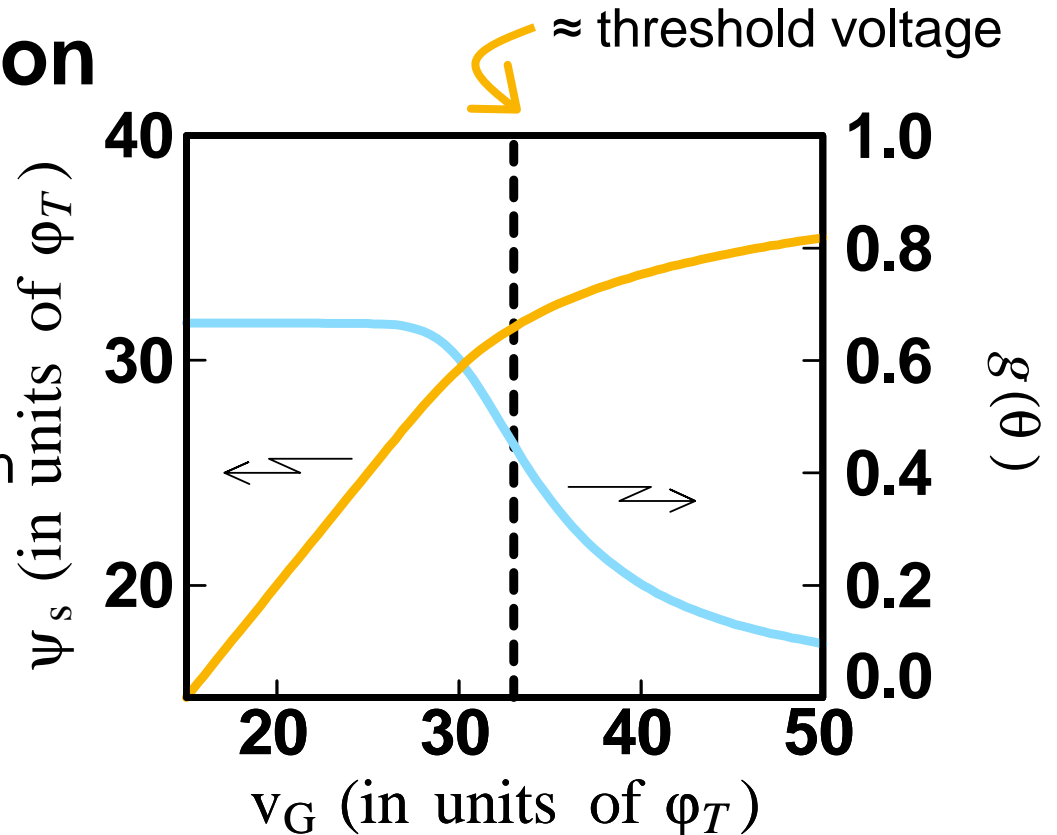
where  $g(\theta) = \frac{\sin(2\theta) - 2\theta \cdot \cos(2\theta)}{\theta \cdot \tan(\theta) \cdot [2\theta + \sin(2\theta)]}$

and  $\frac{\gamma}{4} \approx \frac{1}{12} \cdot \frac{t_{Si}}{t_{ox}} < 1$  (typical value: 1/2)



# new model – derivation

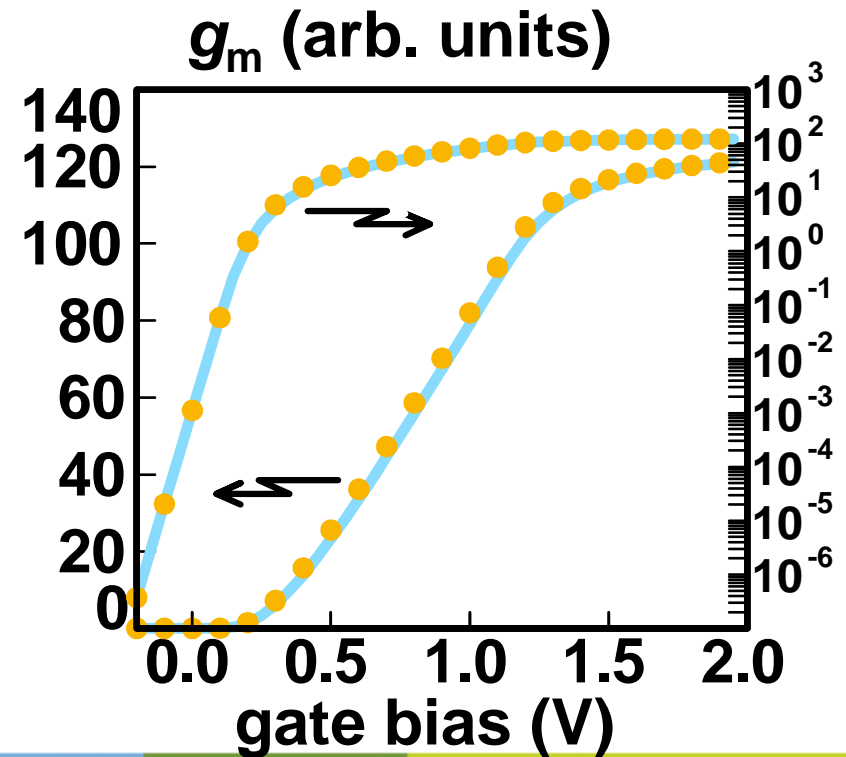
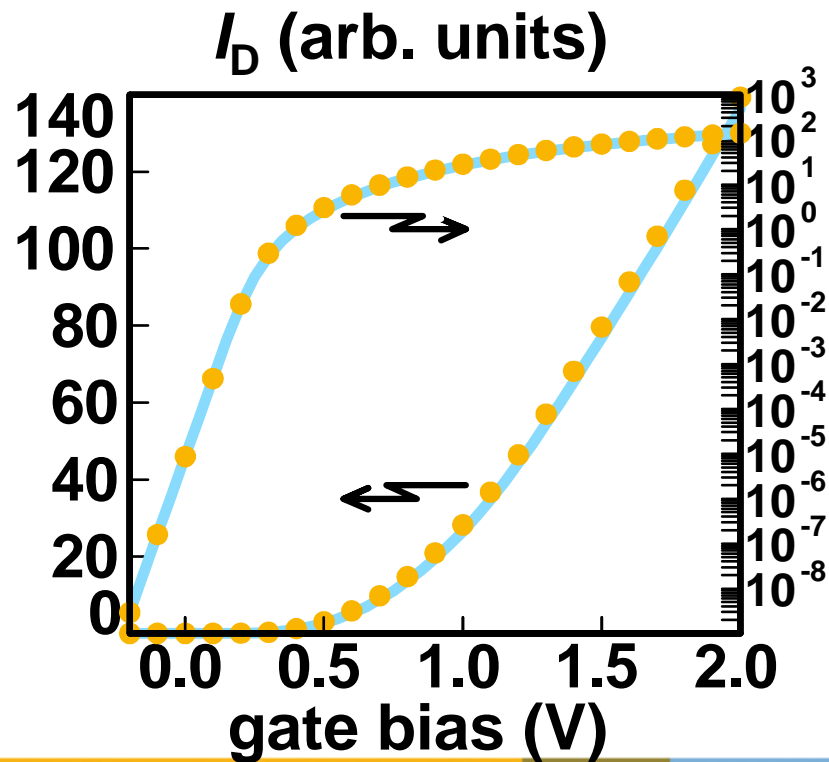
- ▶  $\tilde{Q} = Q \cdot \left[ 1 + \frac{\gamma}{4} \cdot g(\theta) \right]$
- ▶ in strong inversion  $\tilde{Q} \approx Q$
- ▶  $\tilde{Q}$  has *no* influence on diffusion current (subthreshold)
- ▶ **conclusion:** approximation  $\tilde{Q} \approx Q$  valid in all regions of operation
- ▶ **note:** error in *current* is smaller than error in  $Q$ !



$$-\frac{I_D}{\mu \cdot h} = \underbrace{\tilde{Q}(x)}_{\text{exact}} \cdot \frac{d\psi_s}{dx}(x) - \frac{dQ}{dx}(x) \quad \longrightarrow \quad \underbrace{Q(x)}_{\text{approximation}} \cdot \frac{d\psi_s}{dx}(x) - \frac{dQ}{dx}(x)$$

# new model – accuracy

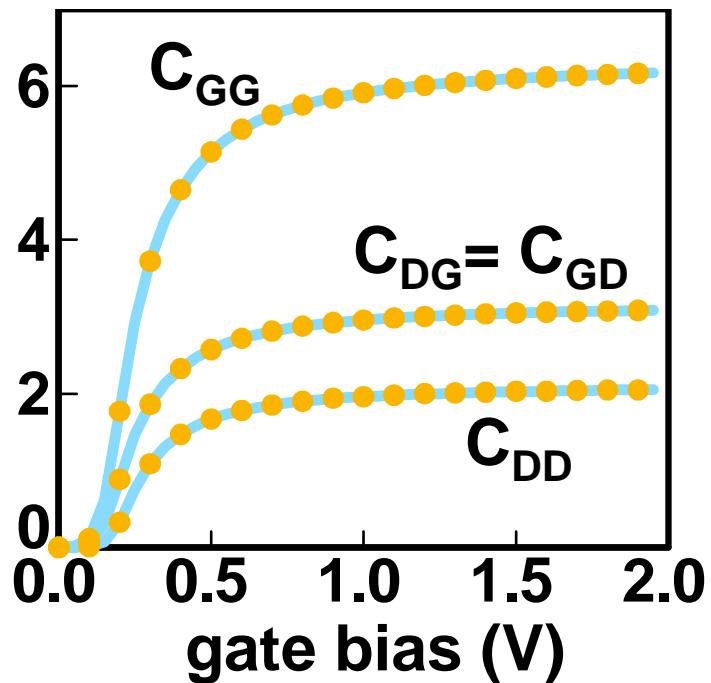
- ▶ high accuracy for currents
  - exact calculation (•) vs. our model (—)
  - $V_{DS}=1V$ , no fitting parameters



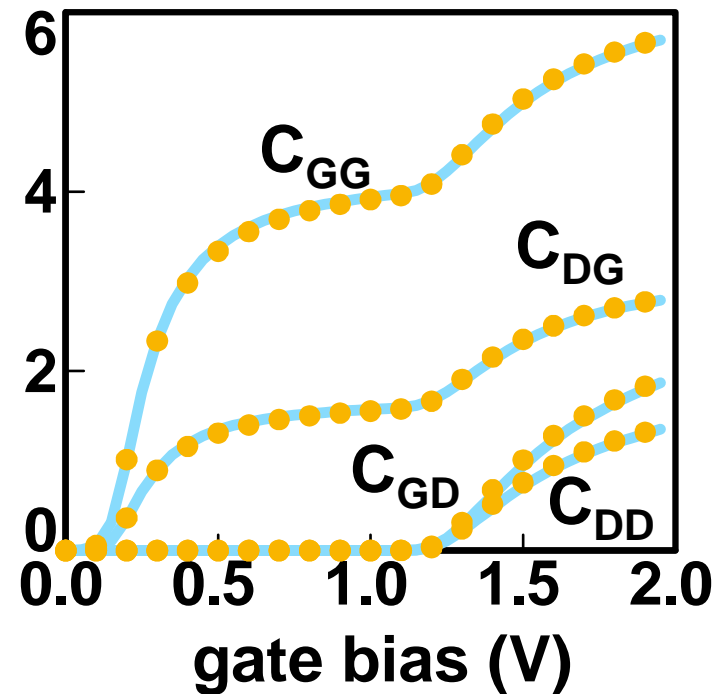
# new model – accuracy

- ▶ high accuracy for charges/capacitances
  - exact calculation (•) vs. our model (—)
  - no fitting parameters

**C (arb. units) @  $V_{DS}=0V$**



**C (arb. units) @  $V_{DS}=1V$**



# new model – equations

► simple expressions

– for charges

$$Q_G = -h \cdot L \cdot \left[ \bar{Q} + \frac{(\Delta Q)^2}{12 \cdot (\bar{Q} - 2 \cdot C_{ox})} \right]$$

$$\bar{Q} = (Q_0 + Q_L) / 2 \quad \Delta Q = Q_L - Q_0$$

$$Q = -\frac{8}{t_{Si}} \theta \tan \theta \quad \theta = \frac{t_{Si}}{4} \sqrt{|\alpha|}$$

– and current

$$I_D = -\frac{\mu \cdot h}{2 \cdot L \cdot C_{ox}} \cdot (\bar{Q} - 2 \cdot C_{ox}) \cdot \Delta Q$$

exact:

$$Q_G \propto \frac{1}{I_D} \left( \left[ \theta^2 - \ln(\cos^2 \theta) - 2\theta \cdot \tan \theta + \frac{4}{3C_{ox} t_{Si}} \theta^3 \tan^3 \theta + h(\theta) \right]_{\theta_0}^{\theta_L} \right)$$

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$$\theta = \frac{t_{Si}}{4} \sqrt{|\alpha|}$$

very suitable for compact model!

# new model – benefits

- ▶ simple expressions for currents and charges (all in terms of  $Q_{x=0}$  and  $Q_{x=L}$ )
- ▶ computationally efficient (no special functions)
- ▶ high accuracy (exact at  $V_{DS}=0$ )
- ▶ approximation preserves essential physics
  - *no* charge-sheet approximation (despite similar equations)
  - volume inversion effect correctly modelled
- ▶ current/charge expressions similar to bulk MOSFET
  - SCEs can be included along the same lines
- ▶ no loss of symmetry properties

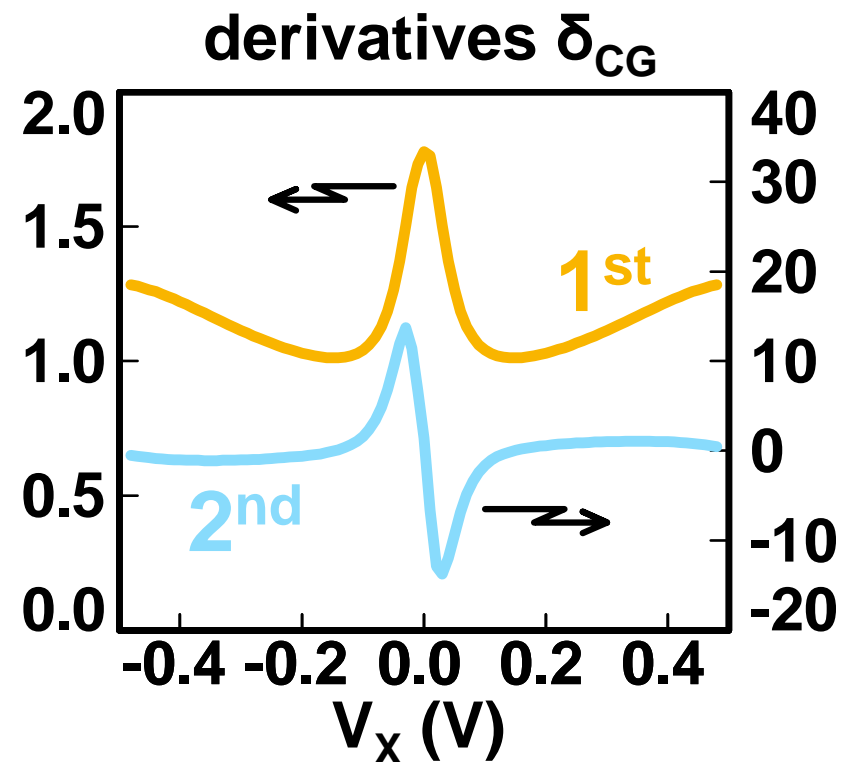
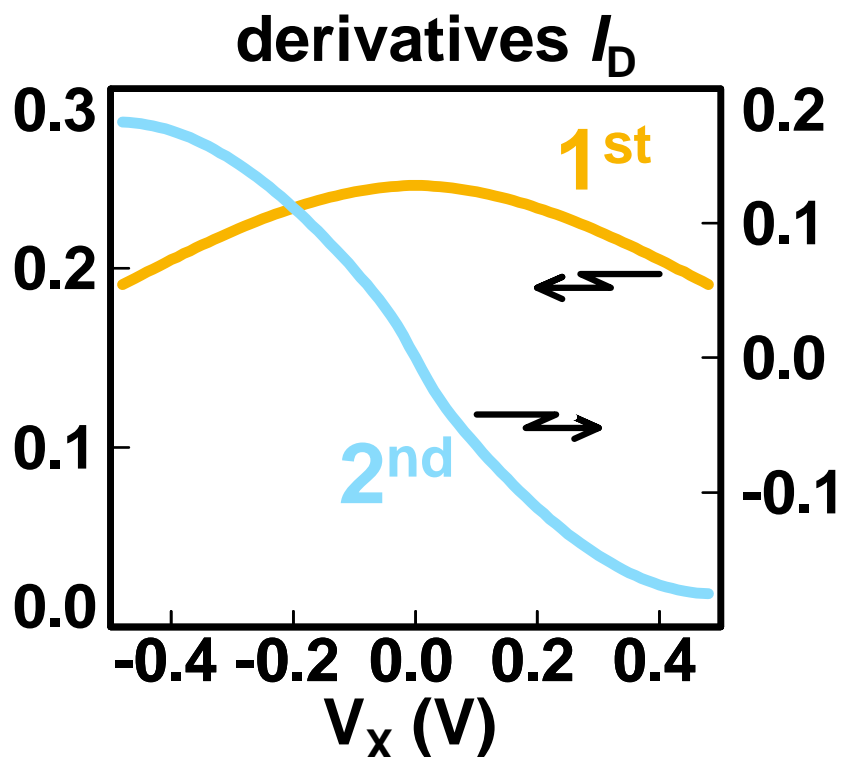


# short channel and additional effects

- ▶ core model describes ideal classical long channel device
- ▶ practical model needs to describe *real* devices, including
  - velocity saturation
  - channel length modulation
  - quantum confinement
  - mobility reduction
  - ...
- ▶ implemented descriptions adopted from PSP bulk-MOSFET model

# Gummel symmetry

- ▶ important for analog simulation
  - related to S/D symmetry
  - see: McAndrew, TED53 p. 2202 (2006)



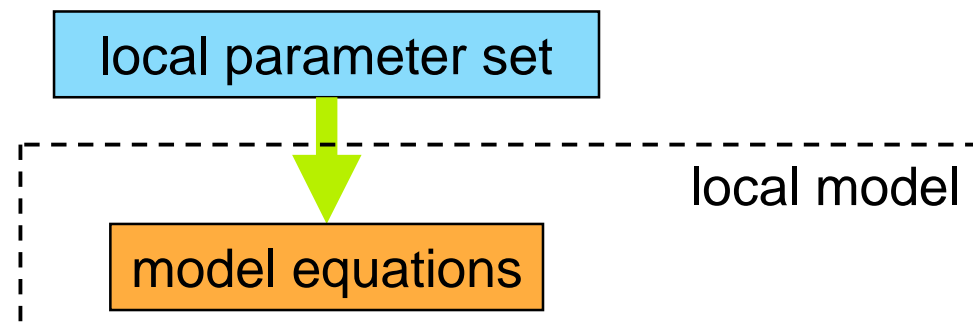
**geometry scaling**

# geometry scaling

- ▶ extraction procedure (similar to PSP)
  - local parameter set for each  $L$
  - fit scaling rules to local parameters
  - create global parameter set
- ▶ 2D-device simulations
  - symmetric FinFET
  - p-channel
  - $L=30\dots400\text{nm}$
  - $t_{\text{Si}}=10\text{nm}$ ,  $t_{\text{ox}}=1.2\text{nm}$

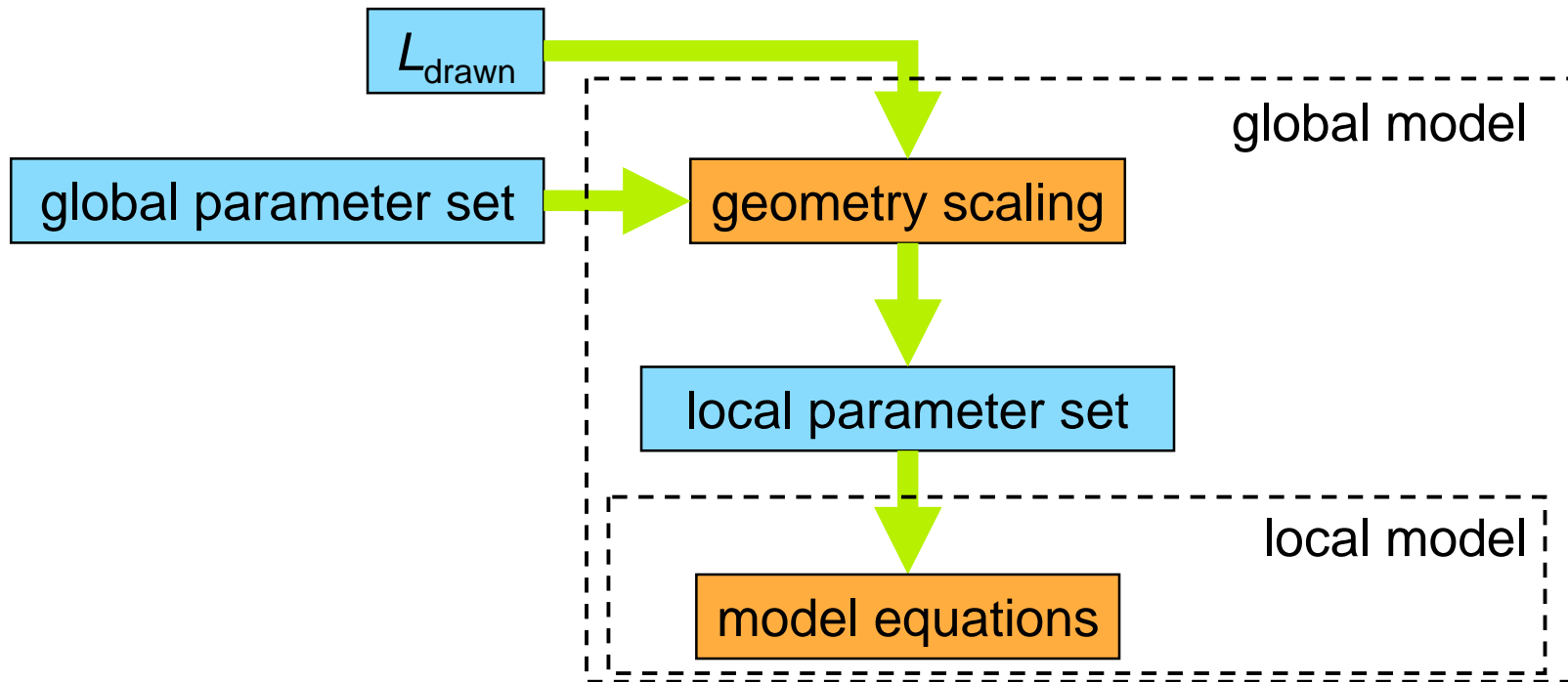
# geometry scaling – local/global model

- ▶ local parameter set: describes one single device (geometry)  
~14 parameters (IV + CV)
- ▶ global parameter set: describes full range of  $L$ -values (~30 parameters)



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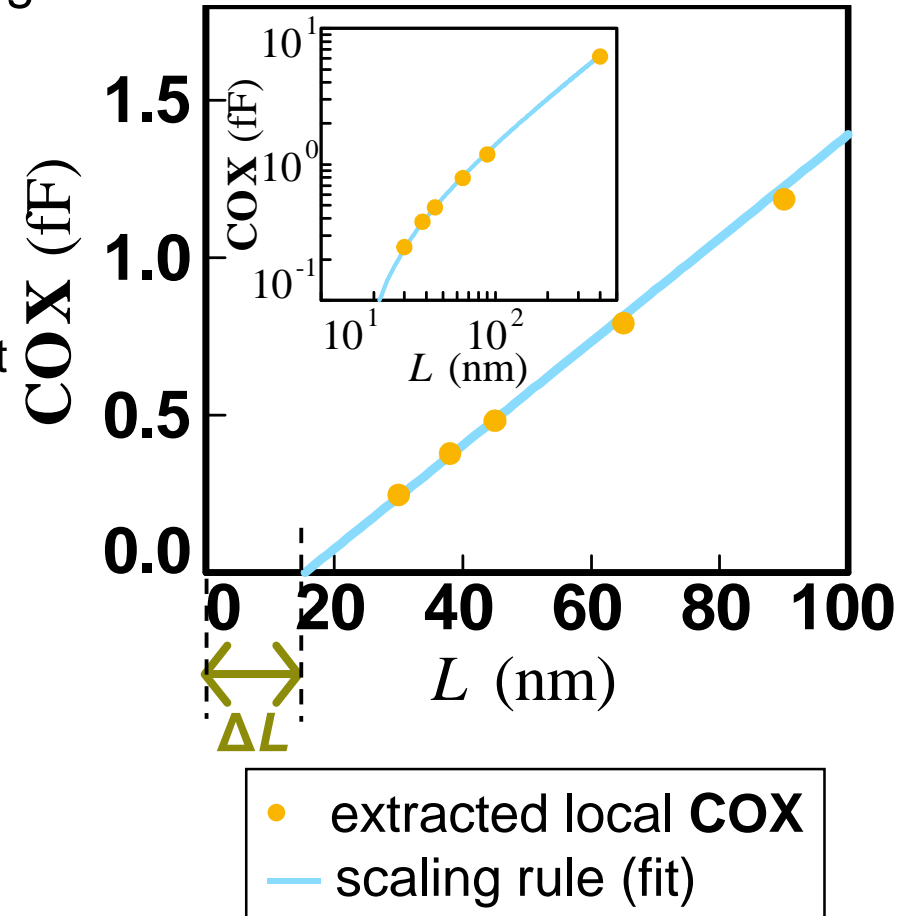
# geometry scaling – effective channel length

- ▶ determination of effective channel length

$$COX = \frac{\epsilon_{ox}}{t_{ox}} \cdot W_{eff} \cdot (L - \Delta L)$$

- ▶ requires accurate value for **COX**
  - usually taken from accumulation capacitance
  - but accumulation capacitance cannot be measured in floating-body device
  - $C_{GG}$  is  $V$ -dependent in inversion
  - QM-effects must be included
- ▶ accurate CV-model required!
- ▶ edge-effects (top gate, buried oxide) included through *effective* fin-height

$$W_{eff} \approx 2 \cdot h_{fin,eff}$$



# geometry scaling – effective channel length

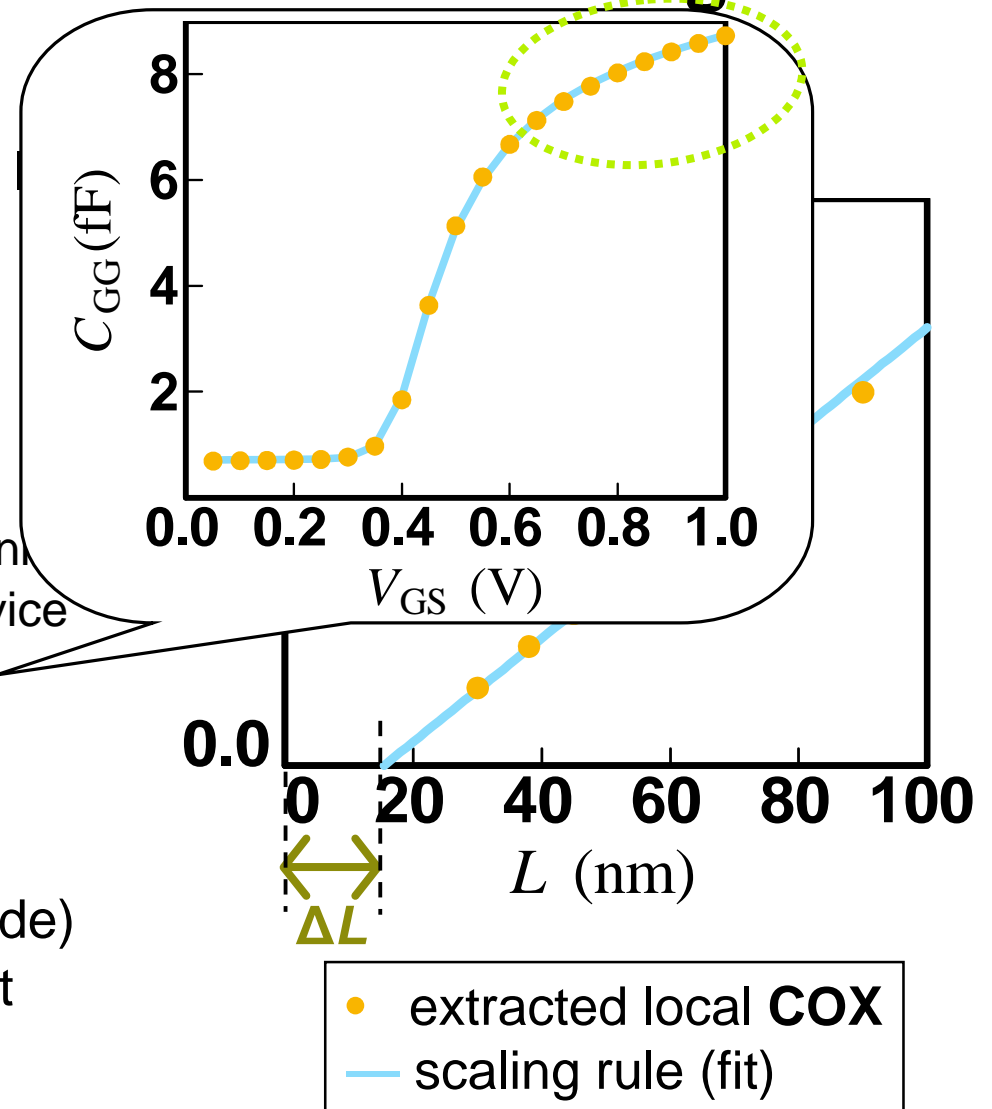
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# geometry scaling – scaling rules

- ▶ fit scaling rules to local model parameters

- ▶ example: non-ideal subthreshold slope  $S$  in short devices

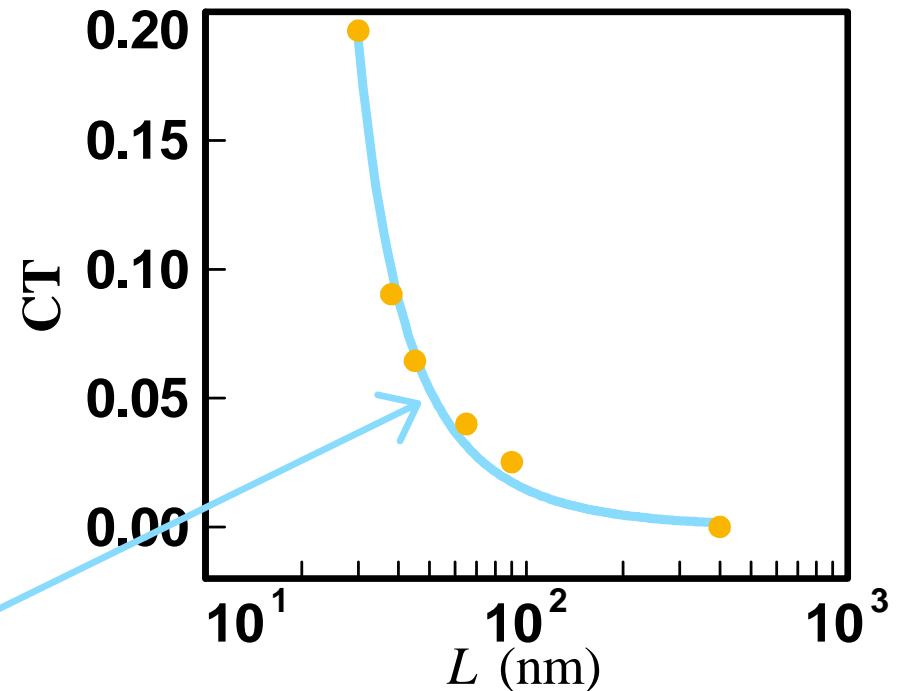
$$S = S_{\text{longchannel}} \cdot (1 + CT)$$

- ▶ scaling rule

$$CT = CTO + CTL \cdot \left( \frac{L_{EN}}{L_E} \right)^{CTLEXP}$$

- ▶ fit to extracted local parameters

$$\begin{aligned} CTO &= 0 \\ CTL &= 0.0004 \\ CTLEXP &= 1.5 \end{aligned}$$



● extracted local CT  
— scaling rule (fit)

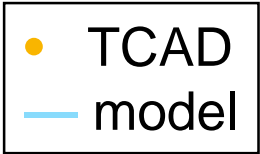
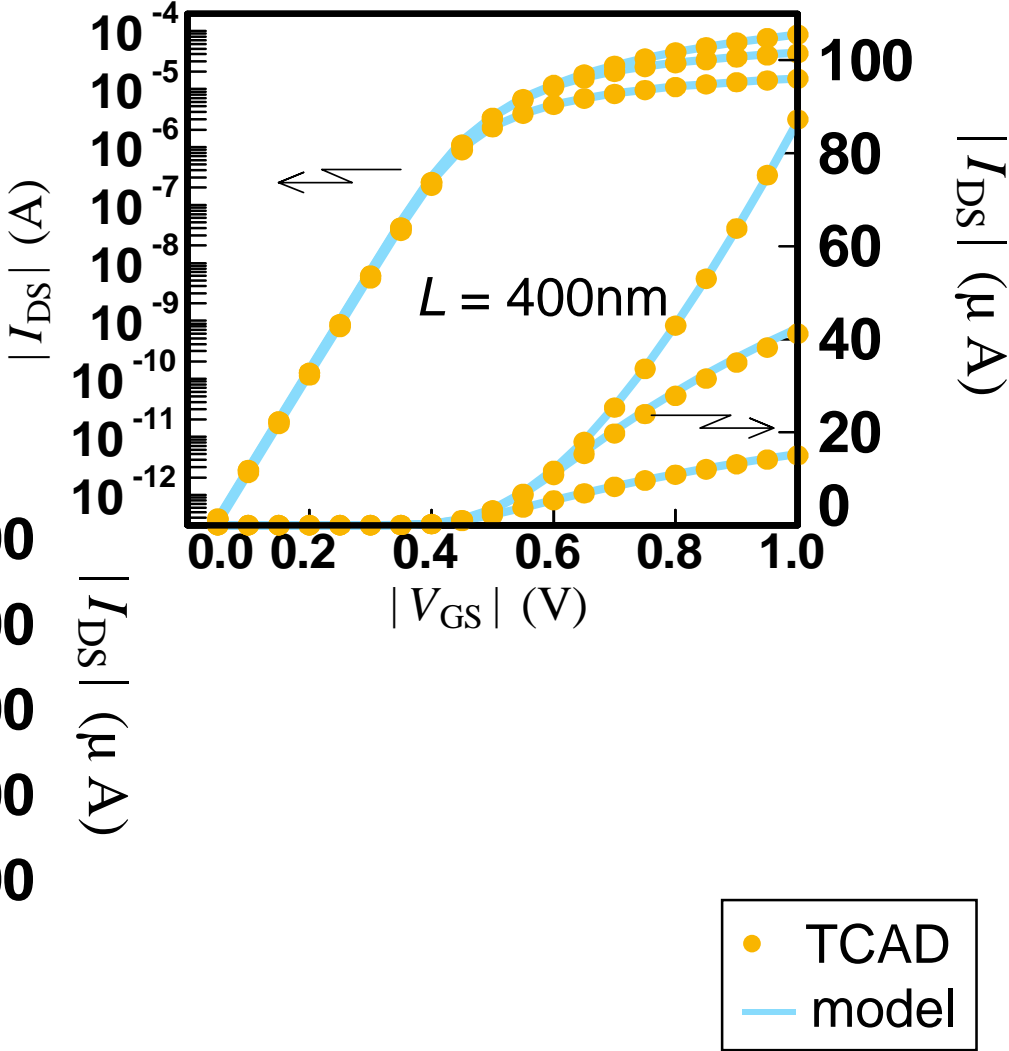
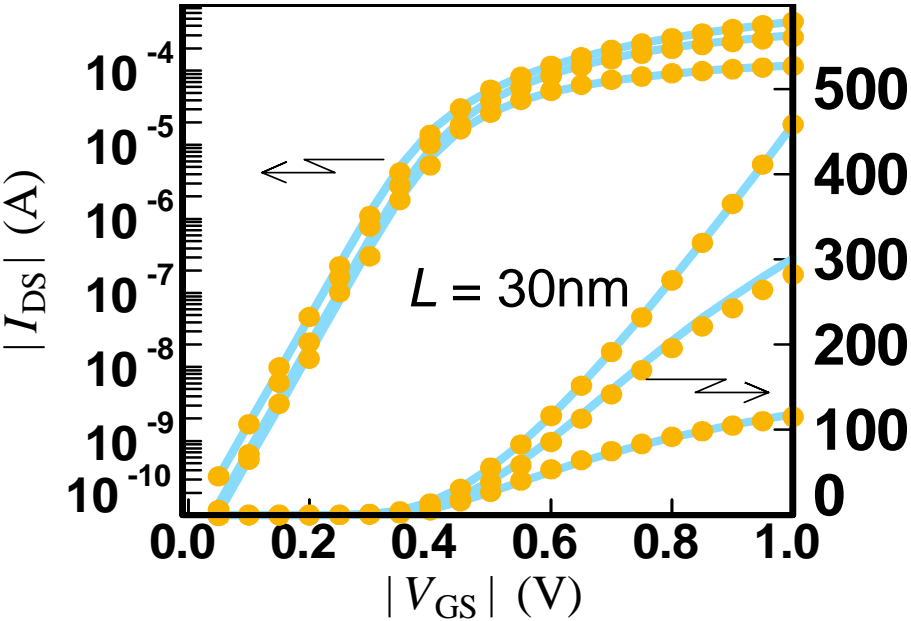
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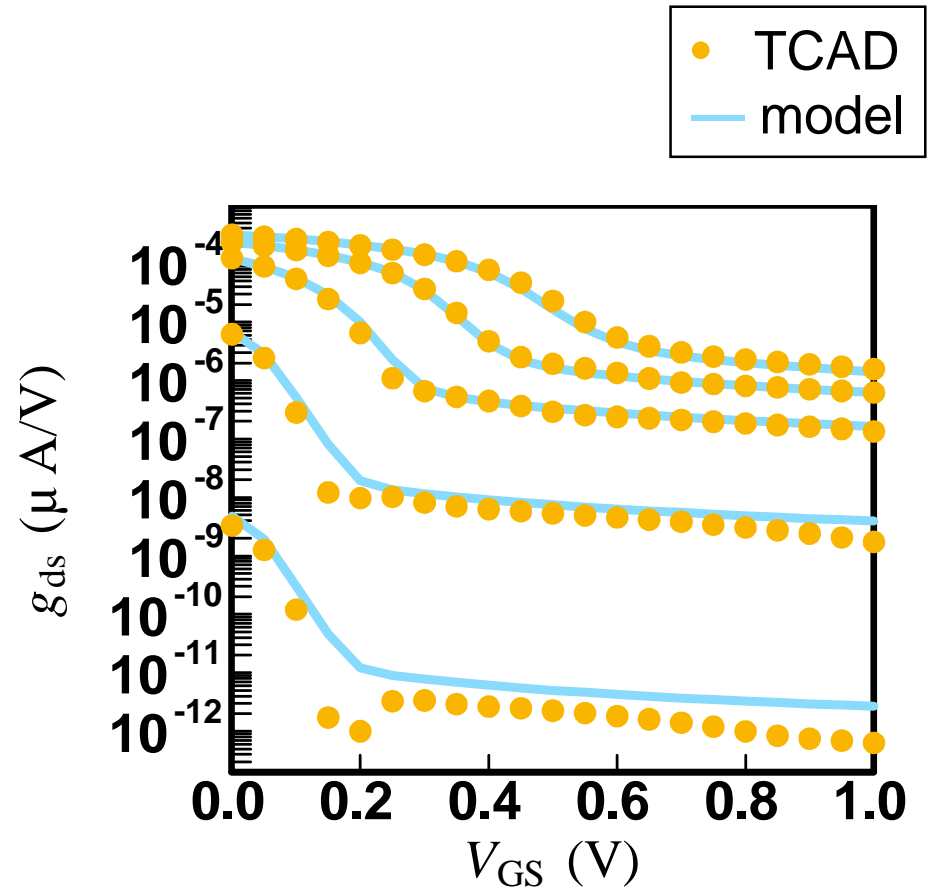
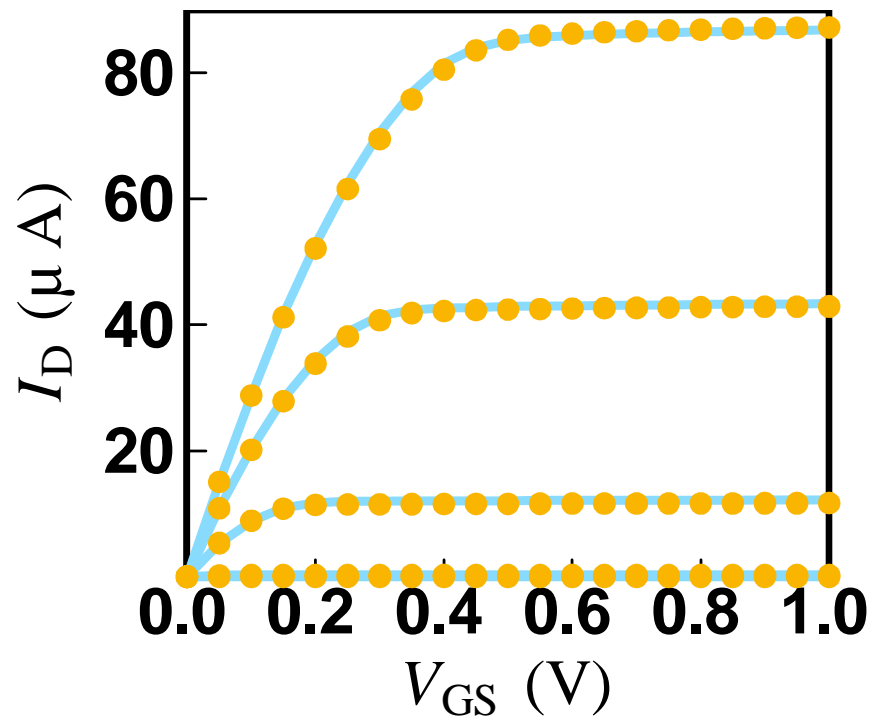
# geometry scaling – results IV

- ▶ Resulting global model fits
- ▶  $V_{DS} = 0.05, 0.15, 1.0 \text{ V}$



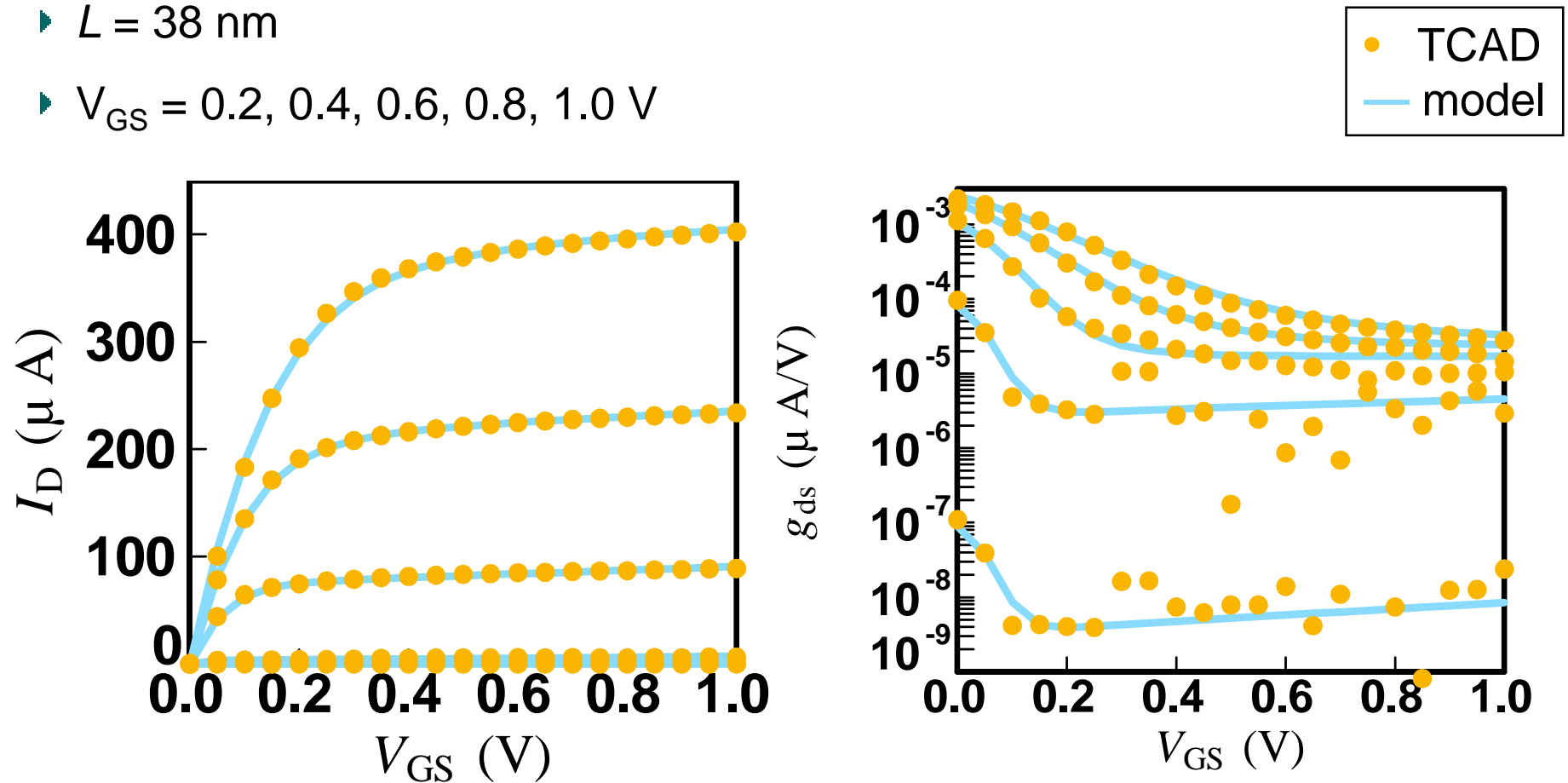
# geometry scaling – results *IV*

- ▶  $L = 400 \text{ nm}$
- ▶  $V_{GS} = 0.2, 0.4, 0.6, 0.8, 1.0 \text{ V}$

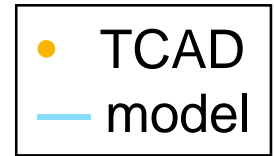


# geometry scaling – results IV

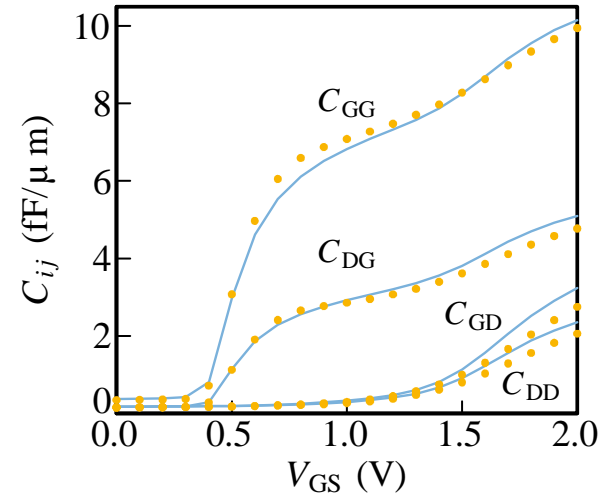
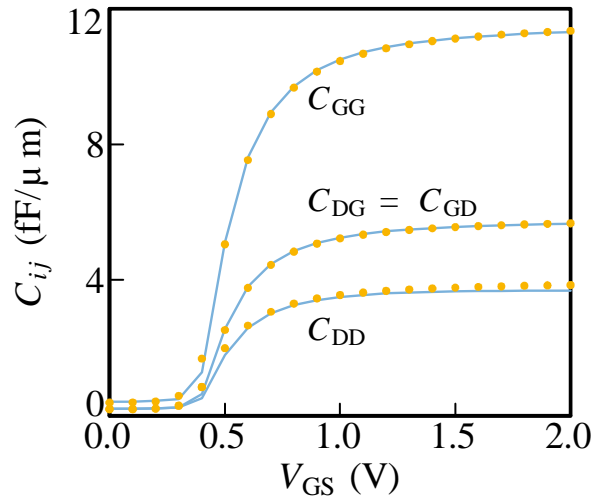
- ▶  $L = 38$  nm
- ▶  $V_{GS} = 0.2, 0.4, 0.6, 0.8, 1.0$  V



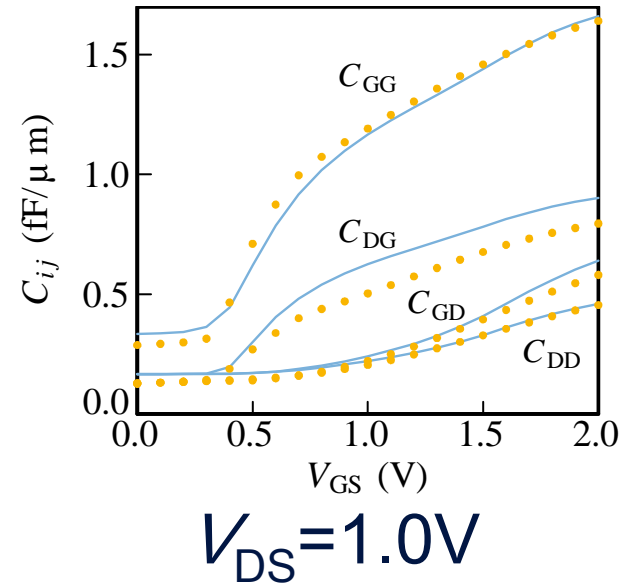
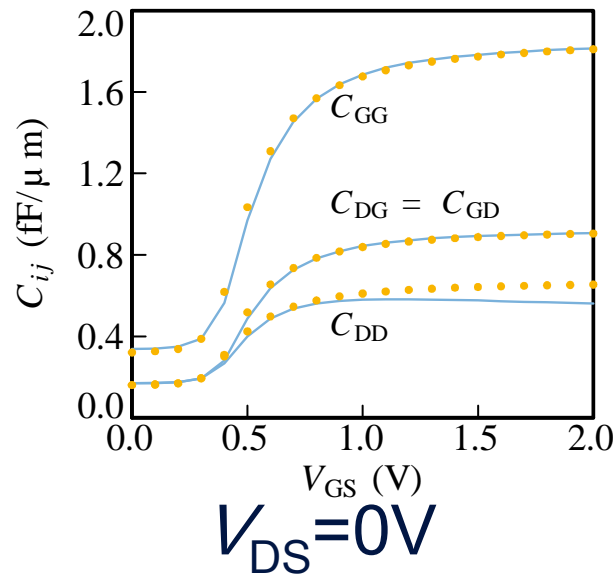
# geometry scaling – results CV



$L=200\text{nm}$

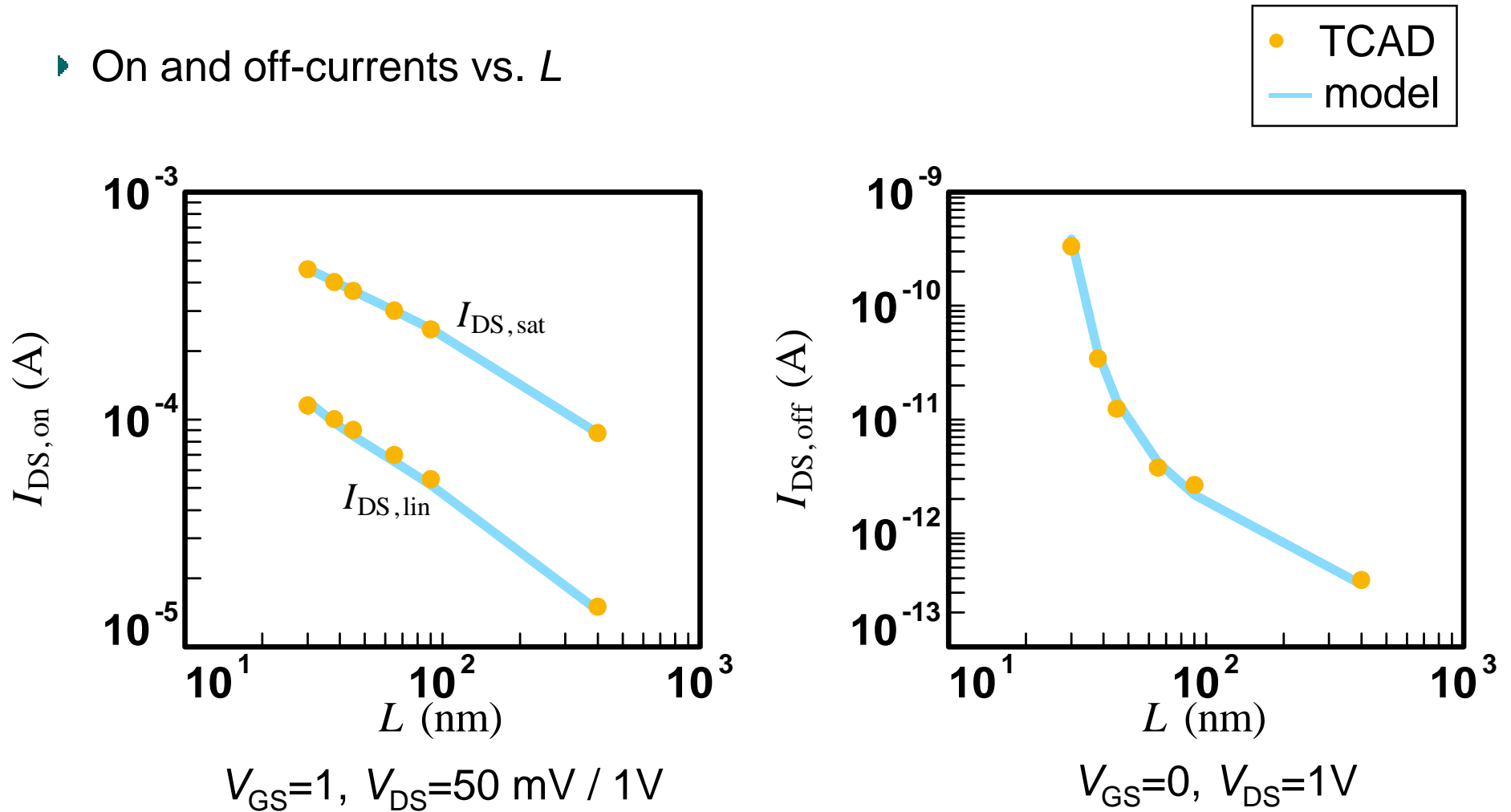


$L=30\text{nm}$



# geometry scaling – on/off currents

- ▶ On and off-currents vs.  $L$



# model overview / summary

- ▶ implemented in Verilog-A
- ▶ true *compact* model
  - no extensive numerical algorithms
- ▶ includes all operating regions
  - sub-threshold
  - linear/triode
  - saturation
- ▶ suitable for all analysis types
  - dc
  - ac
  - transient
  - HB, PSS, ...
- ▶ hierarchy (for easy parameter extraction)
  - local and global model



# conclusion

- ▶ PSP-based compact FinFET model
- ▶ PSP bulk-MOSFET model elements have been reused
- ▶ compact and computationally efficient
- ▶ model gives excellent description of
  - currents (subthreshold, linear, saturation)
  - conductances
  - capacitances
- ▶ model covers full range of channel lengths



